



Chhatrapati Shahu Ji Maharaj
University, Kanpur

Answer Script Details
Barcode 6433500

Roll No. 24077000697
Total Mark 53/75.00

Exam MASTER OF SCIENCE _ODD EXAM-DEC-24
Subject B010702T - CLASSICAL MECHANICS

Question wise Mark Summary

Q.No Mark Q.No Mark Q.No Mark Q.No Mark

1A 4/5 8 NA/15

1B 3/5 9 NA/15

1C 3/5

1D 4/5

1E 3/5

1F 4/5

1G 4/5

1H 4/5

1I 3/5

2 11/15

3 NA/15

4 NA/15

5A NA/5

5B NA/5

5C NA/5

6 10/15

7 NA/15

Chhatrapati Shahu Ji Maharaj University Kanpur, Uttar Pradesh

Date of Exam: 23/01/25 Shift: I Room No: 25
 Paper Code: BD10702T Subject: CLASSICAL MECHANICS Year/Sem: I/I
 Name of Candidate: FARHEEN RAHMAN
 Roll No: 24077000697

Signature of Candidate: *Farheen Rahman*
 Signature of Invigilator: *[Signature]*
 COE Facsimile: *[Signature]*

PART-II

MARKS OBTAINED										
Q.	1	2	3	4	5	6	7	8	9	10
(a)										
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Total										
Total Marks in Figures							Max. Marks			
Total Marks in Words										

B010702T

Paper Code

Signature of Evaluator

Course: M.Sc.
 Session: 2024-25 Year/Semester: I/I
 Subject Name: CLASSICAL MECHANICS
 Medium: English Hindi
 Paper Code: B010702T
 Exam Date: 23012025
 Name of Candidate: FARHEEN RAHMAN
 Father's Name: HABIB UR RAHMAN

संस्थान का कोड
College Code

K	N	0	4
A	A	0	0
E	B	1	1
F	D	2	2
H	J	3	3
K	4	4	4
L	L	5	5
R	M	6	6
S	7	7	7
U	T	8	8
U	9	9	9
W			

परीक्षा केंद्र का कोड
Exam Centre Code

K	N	0	4
A	A	0	0
E	B	1	1
F	D	2	2
H	J	3	3
K	4	4	4
L	L	5	5
R	M	6	6
S	7	7	7
U	T	8	8
U	9	9	9
W			

परीक्षा का प्रकार
Type of Exam

Regular
 Private
 Re-Exam
 Back Paper Exam

ANSWER BOOKLET NO.

6433500

B010702T

Paper Code

Enrollment Number: CSJMA24000130954
 Candidate's Roll Number: 24077000697
 Paper Code: B010702T

2	4	0	7	7	0	0	0	6	9	7
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B	0	1	0	7	0	2	T
A	0	0	0	0	0	0	N
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G	5	5	5	5	5	5	
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7	7	7	7	7	7	7	
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9	9	9	9	9	9	9	

Signature of Candidate: *Farheen Rahman*

Signature of Invigilator: *[Signature]*

C S Facsimile

COE Facsimile: *[Signature]*

नोट - 1. परीक्षार्थी को निर्दिष्ट किया जाता है कि आवरण वाले को पुरत ध्यान पर उचित सभी निर्देशों को सावधानीपूर्वक पढ़ें।
 2. अंकन में भरो जाने वाली प्रतिक्रियाएँ सभी उत्तरों से शुद्ध की जाएँ। 3. नीचे के स्थानों पर नीले अंकन से भरा जाएँ।

INSTRUCTION TO THE CANDIDATE FOR FILLING PART-I

1. Read the instructions carefully given on the answer script and admit card.
2. Write Date of Exam, Shift, Paper Code & Name of Subject Correctly
3. Write Name & Roll No. Correctly.
4. Write Semester & Branch Correctly.

INSTRUCTION TO THE CANDIDATE FOR FILLING PART-III

1. Use blue or black ball point pen for writing alphabets & numerals in boxes.
2. Carefully study the example before you start marking.
3. As shown in the example below, blacken the circles completely.



4. Make no Stray marks on this sheet.

5. DO NOT WRITE OR MARK ON THE BAR CODE.

IN ORDER TO AVOD UFM (UNFAIR MEANS) :

1. The Roll No. and Answer Book no. found elsewhere or any other symbol found in the answer book will be treated as unfair means.
2. Any tempering of Bar Code and Booklet no shall be treated as Unfair Means.
3. Do Not bring the materials like slip of paper/mobile/digital diaries/ study material/ revision notes in examination hall. Possession of the mobiles/ digital diaries/electronic/digital/ watch and any other electronic gadget except memory less scientific calculator shall be considered as UFM case.
4. Do not keep or paste currency note in answer script it shall be consider as UFM.

अनुचित साधन से बचने हेतु :

1. उत्तर पुस्तिका के निर्दिष्ट स्थान को प्रत्येक अनुक्रमिक एवं उत्तरपुस्तिका का क्रमिक कड़ी और न डिले तथा कोई भी चिह्न न बनावे क्योंकि यह अनुचित साधन प्रयोग की परिधि में आता है।
2. उत्तर पुस्तिका के सामनेक अथवा उत्तर पुस्तिका संख्या पर छेद छानद करने पर अनुचित साधन प्रयोग माना जावेगा।
3. परीक्षा कक्ष में निम्न वस्तुएं लाय न लायें, जैसे किताबें हूट कागज के टुकड़ें, मोबाईल, डिजिटल काली, डिजिटल वॉच, काली, प्लासक यह सभी वस्तुएं जो अनुचित साधन को अवलंबित करती है। केवल संबंधित प्रश्नपत्र में ही वेधोरो मेंस स्टूडेंट्सक कोन्सुमेटर ले जाने को अनुमति है।
4. उत्तर पुस्तिकाओं में कपड़े न रखें न ही उत्तर पुस्तिका में लिपिकावे। ऐसा करना अनुचित साधन प्रयोग की परिधि में आता है।

उत्तरपुस्तिकाओं को भरीय भाँटना

1. प्रत्येक पत्र एवं उत्तरपुस्तिका पर दिवसे दिवसे निर्देशों को ध्यान से पढ़ें।
2. उत्तर पुस्तिका के टुकड़ों पर छेद न डिलें।
3. उत्तरपुस्तिका के पृष्ठों पर टोनों लक डिलें।
4. प्रश्न पत्र पर अपने अनुक्रमिक के अतिरिक्त कुछ न डिलें।
5. प्रश्न पत्र कोड एवं प्रश्न पत्र ID सामग्री पूर्णक डिलें।
6. अपनी स्थिति स्पष्ट डिलें।
7. उत्तरपुस्तिका के पृष्ठों की संख्या देखें। उत्तरपुस्तिका में पृष्ठ (1-24) से कम है या कटे हुए है, तो परीक्षा शुरू होने के पूर्व टुकड़ों उत्तरपुस्तिका ले लें।
8. प्रश्नपत्र को देख, यदि प्रश्नपत्र के विषय कोड, विषय का नाम तथा प्रश्न न कोडें मुटि है, तो उसको परीक्षा शुरू होने के 30 मिनट के अन्दर कल निर्देशक को लक्ष्यत सुचित करें, उसके बाद विरक्षिततायव द्वारा कोडें लक नहीं की जावेगी।
9. प्रश्नों के उत्तर लिखने के लिये पेंसिल का प्रयोग न करें।
10. बी कोपी का अतिरिक्त टुक नही दिया जावेगा।

INSTRUCTION TO THE CANDIDATE

1. Read the instructions carefully given on the Question Paper, Admit Card & Answer Script.
2. Do not write anything on back side of the cover page.
3. Write on both sides of pages of answer book.
4. Do not write anything on question paper except Roll Number.
5. Write Paper Code & Question Paper Id carefully.
6. CHECK the number of pages (1-24) or any other kind of damage in your answer script, if found than change the answer script immediately before the commencement of examination.
7. CHECK the Question Paper for any kind of discrepancy e.g. Subject Code, Subject Name, and Question of the Question Paper during first THIRTY MINUTES of the commencement of the exam, so that it can be corrected in TIME. After that no corrections shall be entertained by the university.
8. Do not use pencil for answering the question.
9. Write status correctly e.g. those appearing in carry over papers should fill in status as Carry Over. Those appearing as Ex- Students should fill in status as ex.
10. No supplementary answer book & graph paper will be provided.

INSTRUCTION TO THE CANDIDATE FOR FILLING PART-IV

1. Use blue or black ball point pen for writing alphabets & numerals in Boxes.
2. Use blue or black ball point pen for filling the circles.

	1	8	1	5	4	3	2	1	6	9
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8	8	●	8	8	8	8	8	8	8	8
9	9	9	9	9	9	9	9	9	9	●

Note- If your Roll No. is of 10 digits. Please leave first three columns .



Paper Code

B010702T



1

Section-A

Short Answer Type Questions

Answer no. 1(A)

Principle of Least action \rightarrow Principle of least action is defined as, "the time integral of twice of kinetic energy is called action integral"

$$\int_{t_1}^{t_2} 2T dt = A \quad \text{--- (1)}$$

Definition \rightarrow "The action integral A of twice of kinetic energy is equal to the zero"

$$\int_{t_1}^{t_2} 2T dt = 0 \quad \text{--- (2)}$$

we know that $\dot{q} \left(\frac{\partial T}{\partial \dot{q}} \right) = 2T$

$$\dot{q} \left(\frac{\partial L}{\partial \dot{q}} \right) = 2T$$

$$p \dot{q} = 2T$$

So, equation (2) becomes as

$$\int_{t_1}^{t_2} \sum p \dot{q} dt = 0$$



since

$$\sum \int_{t_1}^{t_2} p \dot{q} dt = 0$$

and

by

Hamilton's principle

$$\delta \int_{t_1}^{t_2} L dt = 0$$

$$H + L = \sum p \dot{q}$$

$$\therefore H + V = T$$

$$H + L = \sum p \dot{q} = 0$$

$$H + L = \text{constant}$$

for the principle of least action.

Hence, principle of least action deduce as

$$\int_{t_1}^{t_2} (2 \times \text{Kinetic energy}) dt = 0$$

$$\int_{t_1}^{t_2} 2T dt = 0$$



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3

Answer : 1(b)

given:

$$L = ax^2 + y^2 - kxy$$

we know $\frac{\partial L}{\partial x} = p_x$ and $\frac{\partial L}{\partial y} = p_y$ (0) : $\frac{\partial L}{\partial x} = 2ax - k$

$$\frac{\partial L}{\partial x} = 2ax - k = 0$$

$$p_x = 2ax - k \Rightarrow x = \frac{p_x + k}{2a}$$

Hamiltonian $H = \sum_i p_i \dot{q}_i - L$

$$H = \sum p \dot{x} - L$$

$$H = p_x \left[\frac{p_x + k}{2a} \right] - [ax^2 + y^2 - kxy]$$

$$H = \frac{p_x^2}{2a} - a \left[\frac{(p_x + k)^2}{4a^2} \right] + \left[\frac{p_y^2}{2} \right] - [kxy]$$

$$H = \frac{p_x^2}{2a} - \frac{p_x^2}{4a} + \frac{p_y^2}{2} - kxy$$

$$H = \frac{p_x^2}{2a} - \frac{p_x^2}{4a} + \frac{p_y^2}{2} - kxy$$

O.M.



Paper Code

B010702T



4

$$H = \frac{1}{4} p_x^2 + \frac{1}{4} p_y^2 - kxy$$

This is the required Hamiltonian.

Answer : 1 (C)

To show: generalised momentum is conserved when corresponding coordinate is cyclic.

If Lagrangian is independent of the coordinate q_k 's then

$$\text{Lagrange: } \frac{\partial L}{\partial q_k} = 0 \quad \text{--- (1)}$$

by Lagrange's equation

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) - \frac{\partial L}{\partial q_k} = 0 \quad \text{--- (2)}$$

by eq. (1)

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) = 0 \quad \text{--- (3)}$$

$$\frac{\partial L}{\partial \dot{q}_k} = \text{constant} \quad \text{--- (4)}$$

We know that $p_k = \frac{\partial L}{\partial \dot{q}_k}$



Paper Code

B0107027



5

so equation (4) becomes,

$$\frac{\partial L}{\partial q_k} = p_k = \text{constant}$$

$$p_k = \text{constant}$$

proved

Here, q_k 's is the cyclic coordinate because due to Lagrange's corresponding momentum is conserved.

Answer: 1 (D)

given; transformation equations

$$Q = q^{\alpha} \cos \beta p \quad \text{--- (1)}$$

$$P = q^{\alpha} \sin \beta p \quad \text{--- (2)}$$

To proof; these equation are canonical

Condition for canonical: equation must satisfied -

$$pdq - PdQ = \epsilon \text{ : exact differential eqn --- (3)}$$

$$dQ = \dot{Q} = d [q^{\alpha} \cos \beta p]$$

$$= q^{\alpha} [-\sin \beta p dp] + \cos \beta p [(\alpha - 1) q^{\alpha-1} dq]$$



eqn (3) becomes.

$$pdq - PdQ = pdq - q^{\alpha} \sin \beta p [-q^{\alpha} \sin \beta p dp + (\alpha q^{\alpha-1} dq)]$$

$$pdq - PdQ = pdq + q^{2\alpha} \sin^2 \beta p dp - \alpha q^{2\alpha-1} \sin \beta p dq$$

$$pdq - PdQ = (p - \alpha q^{2\alpha-1} \sin \beta p) dq + q^{2\alpha} \sin^2 \beta p dp \quad \text{--- (4)}$$

eqn (4) becomes as

$$Mdx - Ndy = 0$$

so $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$ must satisfied

On comparing

$$\frac{\partial}{\partial p} (p - \alpha q^{2\alpha-1} \sin \beta p) = \frac{\partial}{\partial q} (q^{2\alpha} \sin^2 \beta p)$$

$$1 - \alpha q^{2\alpha-1} \cos \beta p = 2\alpha q^{2\alpha-1} \sin^2 \beta p$$

equation becomes eqn. to zero

if $2\alpha - 1 = 0$ $q^0 = 1$



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7

then $\alpha = \frac{1}{2}$ put in last expression
we get
 $\beta = 2$

And given equation is an exact differential equation because it satisfied exact differential equation condition

$$\alpha = \frac{1}{2} \text{ and } \beta = 2 \text{ Ans.}$$

Answer: 1(e)

Inertia Tensor \rightarrow consider the system of n particles having $T_{\alpha\beta}$ matrices and ω is the angular frequency.

It is defined as for normal mode particles oscillate with same frequency and phase and amplitude are not same in normal mode.

It is defined by matrix $T = T^T \omega$

having 3×3 row matrix and contains 9 element



It can be write it as

$$\begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} I_{11} & I_{12} & I_{13} \\ I_{21} & I_{22} & I_{23} \\ I_{31} & I_{32} & I_{33} \end{bmatrix} \begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}$$

which represent the inertia tensor.

Properties of Inertia tensor →

- i) Inertia tensor has \checkmark angular frequency of particles of system oscillation.
- ii) Inertia tensor represent by matrix (in above expression) which relates the angular frequency and tensor quantity.

Answer 3 (f)

Hamilton is the advance part of Lagrange formalism.

In Lagrange has contain generalised coordinate q_i \checkmark



q, \dot{q} and time t which are dependent to each other but Hamilton has generalised coordinates i.e. generalised momentum, generalised position and time which are independent to each other.

and $L = L(q, \dot{q}, t)$

$$H = H(q, p, t)$$

Hamilton coordinates are independent to each other so it is an explicit function of time.

we know $-\frac{\partial L}{\partial t} = \frac{dH}{dt}$

If Lagrange's is independent of time then

$$0 = \frac{dH}{dt}$$

$H = \text{constant}$ proved

Hence, H is not an explicit function of time and H is constant



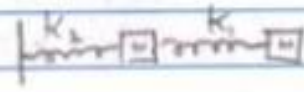
Paper Code

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10

Answer : 1(g)

given the system  is spring
mass system so

$$\text{Kinetic energy } T = \frac{1}{2}(m_1 \dot{x}^2 + m_2 \dot{x}^2)$$

$$\text{and } V = \frac{1}{2} K x^2$$

$$\text{Lagrange } L = T - V$$

$$L = \frac{1}{2} m_1 \dot{x}^2 + \frac{1}{2} m_2 \dot{x}^2 - \frac{1}{2} K x^2 \quad \text{--- (1)}$$

$$\text{by } \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = 0 \quad \text{--- (2)}$$

$$\therefore \frac{\partial L}{\partial x} = 0 - Kx$$

$$\frac{\partial L}{\partial \dot{x}} = m_1 \dot{x} + m_2 \dot{x} \quad \text{put in (2)}$$

$$\frac{d}{dt} [(m_1 + m_2) \dot{x}] - (-Kx) = 0$$

$$(m_1 + m_2) \ddot{x} + Kx = 0$$

$$\ddot{x} = -\frac{Kx}{m_1 + m_2}$$



Paper Code

8010702T



11

$$\text{frequency } \omega = \sqrt{\frac{k}{v}} \quad \text{--- (3)}$$

$$\ddot{x} = \frac{-Rx}{m_1 + m_2}$$

$$k = \frac{(m_1 + m_2)\ddot{x}}{-x} \quad \text{--- (4)}$$

Solve them we get

$$\text{frequency } \omega = \sqrt{\frac{k}{m_1 + m_2}} \quad \text{proved}$$

Answer : 1(h)

Principle of virtual work \rightarrow Let a motion imposed on the body in static equilibrium. which we represent virtual situation. the individual force is zero so the complete work done by the body is also zero.

This is done, Work done by the force of constrain as way that virtual displacement is zero.

$$\delta W_i = \sum F_i (\delta x_1, \delta x_2, \dots, \delta x_n, \dots, \delta x_{f+1}, \dots) \quad \text{--- (5)}$$

$$W = F \times \text{Disp.}$$

$$\delta W_i = \sum F_i \times \delta x_i \quad \text{--- (2)}$$



for total force,

$$F_i = F_i^a + f_i \quad \text{--- (3)}$$

in equilibrium state, work done becomes

$$\delta W_i = \sum_j (F_j^a + f_j) \cdot \delta x_j = 0 \quad \text{--- (4)}$$

here force of constraints $f_i = 0$

becomes as

$$\delta W_i = \sum_j (F_j^a) \cdot \delta x_j = 0$$

Ans. \leftarrow virtual work done by all particles

$$\delta W_i = \sum_j F_j \cdot \delta x_j = 0$$

This is the required principle of virtual work.

Answer : 1(i)

Theory of small oscillation \rightarrow Theory of small oscillation is for large \checkmark particle system having i^{th} particle of the given system of oscillation.



Normal modes \rightarrow If system of particle having same frequency of oscillation then called it as normal mode.

Normal mode does \checkmark exist in same phase and same amplitude of system.

Oscillation of i^{th} particle of system are represented by

$$|\dot{V} - \omega \dot{T}| = 0$$

where \dot{V} = function of Potential Energy magnitude

\dot{T} = Magnitude of Kinetic Energy

ω = angular frequency of i^{th} Particle

\dot{q}_1, \dot{q}_2 = function of V_1 and V_2

\dot{q}_1^2 = function of V_1 \checkmark

\dot{q}_2^2 = function of V_2

This theory explain the system of particles having same frequency of oscillation for system.



Section - B

Answer: 2

Generalised coordinate \rightarrow The minimum numbers of independent coordinates that describe the motion of particle of system compatible with constraints known as 'generalised coordinate'.

transformation equation for generalised coordinates are

$$\left. \begin{aligned} x_i &= x_i(q_1, q_2, \dots, q_f, \dots, q_n, t) \\ y_i &= y_i(q_1, q_2, \dots, q_f, \dots, q_n, t) \\ z_i &= z_i(q_1, q_2, \dots, q_f, \dots, q_n, t) \end{aligned} \right\} - 1(a)$$

for position vector

$$\vec{r}_i = \vec{r}_i(q_1, q_2, \dots, q_f, \dots, q_n, t) - 1(b)$$

1(a) and 1(b) are transformation eqn

D'Alembert principle for generalised coordinate \rightarrow

we can derive the Lagrange equation



of motion by D'Alembert principle, which defined as "The resultant force together with the reverse effective force is in equilibrium".

by Newton's 2nd law of motion

$$F = \frac{dp}{dt} = \dot{p} \quad \text{---(1)}$$

virtual displacement $\frac{\partial x_i}{\partial q_k} \delta q_k = \delta x_i$

and work done

$$\delta W = F \times \delta x_i$$

$$\sum_i (F - \dot{p}) \delta x_i = 0 \quad \text{---(2)}$$

for total force $F = F_i^a + f_i \quad \text{---(3)}$

$$\sum_i (F_i^a + f_i) \delta x_i - \sum_i \dot{p}_i \cdot \delta x_i = 0$$

here force of constraint are zero

$$\sum_i (F_i^a + 0) \delta x_i - \sum_i \dot{p}_i \cdot \delta x_i = 0$$

write it as

$$\sum_i (F_i - \dot{p}_i) \cdot \delta x_i = 0$$

This is the D'Alembert principle.



for generalised coordinate q_k

$$(F_k - \dot{p}_k) \cdot \delta q_k = 0$$

Lagrange equation of motion \rightarrow

We can derive Lagrange's equation of motion from D'Alembert principle for taking system of i^{th} particle having position vector

$$\vec{r}_i = \vec{r}_i(q_1, q_2, \dots, q_n, t) \quad \text{--- (1)}$$

and

$$\text{virtual displacement is } \sum_i \vec{F}_i \cdot \delta \vec{r}_i = 0 \quad \text{--- (2)}$$

$$\delta \vec{r}_i = \frac{\partial \vec{r}_i}{\partial q_k} \delta q_k + \frac{\partial \vec{r}_i}{\partial t} \delta t$$

By definition of D'Alembert principle,

$$(\vec{F}_i - \dot{\vec{p}}_i) \cdot \delta \vec{r}_i = 0$$

$$\vec{F}_i \cdot \delta \vec{r}_i = \dot{\vec{p}}_i \cdot \delta \vec{r}_i \quad \text{--- (3)}$$

Take LHS

$$\vec{F}_i \cdot \delta \vec{r}_i = m_i \sum_k \vec{F}_i \cdot \frac{\partial \vec{r}_i}{\partial q_k} \delta q_k$$

$$\text{Let } G_k = \vec{F}_i \cdot \frac{\partial \vec{r}_i}{\partial q_k}$$

$$\vec{F}_i \cdot \delta \vec{r}_i = \sum_k G_k \delta q_k \quad \text{--- (4)}$$



Take LHS

$$\vec{p}_i \cdot \vec{v}_i = \sum m_i \vec{v}_i$$

$$= \sum \frac{d}{dt} (m_i) \vec{v}_i \quad \text{--- (5)}$$

$$\therefore \left[\frac{d}{dt} (m_i \vec{v}_i) = m_i \frac{d}{dt} (\vec{v}_i) + \vec{v}_i \frac{d}{dt} (m_i) \right]$$

put in above eqn (5)

$$= \sum \frac{d}{dt} (m_i \vec{v}_i) - m_i \frac{d}{dt} (\vec{v}_i)$$

$$= \sum \frac{d}{dt} (m_i \vec{v}_i) - m_i \frac{d}{dt} \vec{v}_i$$

$$= \sum \frac{d}{dt} \left(\frac{m_i v_i^2}{2} - \frac{1}{2} m_i v_i^2 \right)$$

$$\therefore T = \frac{1}{2} m v^2$$

$$= \left[\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}} \right) - \frac{\partial T}{\partial q} \right] \quad \text{--- (6)}$$

now L.H.S = R.H.S

$$\boxed{\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}} \right) - \frac{\partial T}{\partial q} = G_{1k}} \quad \text{--- (7)}$$

This is the general form of Lagrange's equation of motion.



For Conservative System \rightarrow force is derived by some potential

$$Q_{ix} = -\frac{\partial V}{\partial q_i} \quad \text{--- (a)}$$

we know

$$L = T - V \quad \text{--- (b)}$$

(a) and (b) put in eq. (7)

$$\frac{d}{dt} \left[\left(\frac{\partial L}{\partial \dot{q}_i} \right) + \frac{\partial V}{\partial \dot{q}_i} \right] - \frac{\partial L}{\partial q_i} + \frac{\partial V}{\partial q_i} = Q_{ix} \quad \text{--- (c)}$$

use $Q_{ix} = -\frac{\partial V}{\partial q_i}$

And Potential energy is not a function of velocity for conservative system.

for conservative system, $\frac{\partial V}{\partial \dot{q}_i} = 0$

then eq. (c) becomes

$$\boxed{\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0}$$

This is the Lagrange equation of motion for conservative system, where L is Lagrangian.



Section - C

Hamilton equation of motion

Let a particle of mass m is moving with velocity v about force f with mass m directed towards centre.

then particle contain energies like potential and kinetic energy.

Hamilton eqⁿ in Cartesian Coordinate

from Cartesian system
particle has kinetic energy

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

and

Potential energy

$$V = V(x, y, z)$$

so

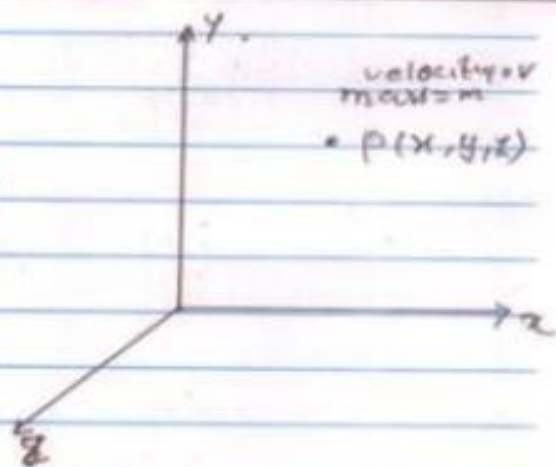
Lagrangian $L = T - V$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(x, y, z)$$

generalised momenta -

$$p_x = \frac{\partial L}{\partial \dot{x}} = \frac{1}{2} \cdot 2m\dot{x} = m\dot{x} \quad \text{or} \quad \dot{x} = \frac{p_x}{m}$$

$$p_y = \frac{\partial L}{\partial \dot{y}} = m\dot{y} \quad \Rightarrow \quad \dot{y} = \frac{p_y}{m}$$





$$p_x = \frac{\partial L}{\partial \dot{x}} = m\dot{x} \Rightarrow \dot{x} = \frac{p_x}{m}$$

Hamiltonian $H = T + V(x, y, z)$

$$H = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + V(x, y, z)$$

$$H = \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m} + V(x, y, z)$$

Hamilton equations are

$$\frac{\partial H}{\partial p_x} = \dot{x} \Rightarrow \frac{p_x}{m} = \dot{x}$$

$$\frac{\partial H}{\partial p_y} = \dot{y} \Rightarrow \frac{p_y}{m} = \dot{y}$$

$$\frac{\partial H}{\partial p_z} = \dot{z} \Rightarrow \frac{p_z}{m} = \dot{z}$$

and

$$\frac{\partial H}{\partial x} = -\dot{p}_x \Rightarrow -\frac{\partial V}{\partial x} = \frac{\partial H}{\partial x}$$

$$\frac{\partial H}{\partial y} = -\dot{p}_y \Rightarrow -\frac{\partial V}{\partial y} = \frac{\partial H}{\partial y}$$

$$\frac{\partial H}{\partial z} = -\dot{p}_z \Rightarrow -\frac{\partial V}{\partial z} = \frac{\partial H}{\partial z}$$



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these are required Hamilton eqⁿ for Cartesian coordinate

Hamilton eqⁿ in cylindrical coordinate

Kinetic energy $T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$

for cylindrical (r, θ, z)

$$V = V(r, \theta, z)$$

$$L = T - V$$

Lagrange $L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(r, \theta, z)$

generalised momenta,

$$p_x = \frac{\partial L}{\partial \dot{x}} \Rightarrow m\dot{x}$$

$$p_\theta = \frac{\partial L}{\partial \dot{\theta}} \Rightarrow m r^2 \dot{\theta}$$

$$H = T + V$$

$$H = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + V(r, \theta, z)$$

or

$$\text{Hamilton } H = \frac{p_x^2}{2m} + \frac{p_\theta^2}{m r^2} + V(r, \theta, z)$$



Canonical equations are

$$\frac{\partial H}{\partial p_j} = \dot{q}_j$$

$$\frac{\partial H}{\partial q_j} = -\dot{p}_j$$



$$\frac{\partial H}{\partial p_j} = \dot{q}_j \Rightarrow$$

$$\boxed{\frac{p_j}{m} = \dot{q}_j} \quad - (1)$$

$$\frac{\partial H}{\partial p_\theta} = \dot{\theta} \Rightarrow$$

$$\boxed{\frac{p_\theta}{r m \dot{\theta}} = \dot{\theta}} \quad - (2)$$

~~$\frac{\partial H}{\partial \theta}$~~ and

$$\frac{\partial H}{\partial \theta} = -\dot{p}_\theta \Rightarrow$$

$$0 \text{ (since } \dot{\theta} \text{ is constant)} + \frac{\partial V}{\partial \theta} = -\dot{p}_\theta \quad - (3)$$

$$\boxed{\frac{\partial H}{\partial \theta} = -\dot{p}_\theta = \frac{\partial V}{\partial \theta}} \quad - (4)$$

These are the Hamilton equations in cylindrical.





Hamilton eqⁿ in spherical coordinate

$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2 + r^2 \sin^2 \theta \dot{\phi}^2) - V(r, \theta, \phi)$$

momenta

$$p_r = \frac{\partial L}{\partial \dot{r}} \Rightarrow m \dot{r}$$

$$p_\theta = \frac{\partial L}{\partial \dot{\theta}} \Rightarrow m r^2 \dot{\theta}$$

$$p_\phi = \frac{\partial L}{\partial \dot{\phi}} \Rightarrow m r^2 \sin^2 \theta \dot{\phi}$$

$$H = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2 + r^2 \sin^2 \theta \dot{\phi}^2) + V(r, \theta, \phi)$$

or

$$H = \left(\frac{p_r^2}{m} + \frac{p_\theta^2}{m r^2} + \frac{r^2 \sin^2 \theta p_\phi^2}{m^2 r^4 \sin^4 \theta} \right) + V(r, \theta, \phi)$$

or

$$H = \frac{p_r^2}{m} + \frac{p_\theta^2}{2 m r^2} + \frac{p_\phi^2}{m^2 r^2 \sin^2 \theta}$$

Hamilton eqⁿ are

$$\frac{\partial H}{\partial p_r} = \frac{p_r}{m} \quad \text{--- (1)}$$

$$\frac{\partial H}{\partial p_\theta} = \frac{p_\theta}{m r^2} \quad \text{--- (2)}$$

$$\frac{\partial H}{\partial p_\phi} = \frac{p_\phi}{m^2 r^2 \sin^2 \theta} \quad \text{--- (3)}$$



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$$\frac{\partial H}{\partial x} = -\dot{p}_x \Rightarrow \frac{\partial V}{\partial x} \quad \text{--- (4)}$$

$$\frac{\partial H}{\partial \theta} = -\dot{p}_\theta \Rightarrow \frac{\partial V}{\partial \theta} \quad \text{--- (5)}$$

$$\frac{\partial H}{\partial \phi} = -\dot{p}_\phi \Rightarrow \frac{\partial V}{\partial \phi} \quad \text{--- (6)}$$

These are Hamilton eqn for spherical system

X

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